Soft photon QED corrections to $B \rightarrow K\ell^+\ell^-$

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Introduction Matrix elements Objectives • $\mathcal{O}(\alpha)$ QED corrections Results Summary and conclusion



Flavour changing neutral currents (FCNCs) are both loop and CKM suppressed

 \implies important candidates to test the Standard Model

Introduction

- Flavour changing neutral currents (FCNCs) are both loop and CKM suppressed \implies important candidates to test the Standard Model
 - $R_{K}^{\mu e} \equiv \frac{\int_{1GeV^{2}}^{6GeV^{2}} dq^{2} \frac{d\Gamma(B^{0} \to K^{0}\mu^{+}\mu^{-})}{dq^{2}}}{\int_{1GeV^{2}}^{6GeV^{2}} dq^{2} \frac{d\Gamma(B^{0} \to K^{0}e^{+}e^{-})}{dq^{2}}} \qquad R_{K}^{\mu e}|_{SM} = 1.00 \pm 0.01 \qquad R_{K}^{\mu e}|_{exp} = 0.846^{+0.060+0.016}_{-0.054-0.014}$

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• The effective Hamiltonian for $b \to s \ell^+ \ell^-$ transition: $H_{eff} = 4 \frac{G_F}{\sqrt{2}} V_{ts}^* V_{tb} \sum_{i=1}^{10} C_i(\mu) \mathcal{O}_i(\mu)$

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• Most relevant operators: $\mathcal{O}_7, \mathcal{O}_9 \& \mathcal{O}_{10}$ • Matrix element for non-radiative decay:

$$M_0(B \to K l^+ l^-) = \frac{G_F \alpha}{\sqrt{2}} V_{ts}^* V_{tb} \left[\left(\begin{cases} C_0^{eff} f_+ + C_7^{eff} \frac{2f_T m_b}{m_{tb}} \end{cases} \right) p^{\mu} - \frac{1}{2} f_{ts} V_{tb} \right]$$



 $\frac{G_F \alpha}{2\sqrt{2\pi}} V_{ts}^* V_{tb} \left[\left(\left\{ C_9^{eff} f_+ + C_7^{eff} \frac{2f_T m_b}{m_B + m_k} \right\} p^{\mu} + \left\{ C_9^{eff} f_- - C_7^{eff} \frac{2f_T m_b}{q^2} (m_B - m_k) \right\} q^{\mu} \right) \left(\bar{l} \gamma_{\mu} l \right) - \left(C_{10} f_+ p^{\mu} + C_{10} f_- q^{\mu} \right) \left(\bar{l} \gamma_{\mu} \gamma_5 l \right) \right] - \left(C_{10} f_+ p^{\mu} + C_{10} f_- q^{\mu} \right) \left(\bar{l} \gamma_{\mu} \gamma_5 l \right) \right] - \left(C_{10} f_+ p^{\mu} + C_{10} f_- q^{\mu} \right) \left(\bar{l} \gamma_{\mu} \gamma_5 l \right) \right]$



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transition:
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$$+\left\{C_{9}^{eff}f_{-}-C_{7}^{eff}\frac{2f_{T}m_{b}}{q^{2}}(m_{B}-m_{k})\right\}q^{\mu}\right)\left(\bar{l}\gamma_{\mu}l\right)-\left(C_{10}f_{+}p^{\mu}+C_{10}f_{-}q^{\mu}\right)\left(\bar{l}\gamma_{\mu}l\right)$$

$$(p_{1}, p_{1}) + e\epsilon_{\alpha}(k)\bar{u}(p_{2})\gamma^{\alpha} \frac{(\gamma_{\mu}p_{2}^{\mu} + \gamma_{\mu}k^{\mu}) + m_{l}}{2p_{2} \cdot k} \Gamma_{A}^{\mu}v(p_{3}) \otimes H_{A\mu}(p_{0}, p_{1})$$

$$\frac{2p_{1}^{\alpha}}{2p_{1} \cdot k}\bar{u}(p_{2})\Gamma_{A}^{\mu}v(p_{3}) \otimes H_{A\mu}(p_{0}, p_{1} + k)$$



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$$p_{1} + e\epsilon_{\alpha}(k)\bar{u}(p_{2})\gamma^{\alpha}\frac{(\gamma_{\mu}p_{2}^{\mu} + \gamma_{\mu}k^{\mu}) + m_{l}}{2p_{2} \cdot k}\Gamma_{A}^{\mu}v(p_{3}) \otimes H_{A\mu}(p_{0}, p_{1})$$

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Leptonic part	Gauge invariant
Hadronic part	Not gauge invariant
Total amplitude	Not gauge invariant







- To fix the gauge invariance of matrix element.
- To get the $\mathcal{O}(\alpha)$ QED correction for the decay width $(B \to K \ell^+ \ell^-)$ and $R_K^{\mu e}$.
- To discuss the collinear divergences and their cancellation.



Real photon emission

$O(\alpha)$ QED corrections

• The addition of a CT $\left(e(Q_B + Q_K)\xi_A k_\mu \left[\bar{u}(p_2)\Gamma^\mu_A v(p_3)\right]\right)$ is required to preserve gauge invariance



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- The non-IR term is important to see the cancellation of collinear divergences.
- Charge conservation and integrating over photons momentum k gives:

$$\tilde{B}_{ij} = \frac{Q_i Q_j \eta_i \eta_j}{2\pi} \left\{ \ln\left(\frac{k_{max}^2 m_i m_j}{\lambda^2 E_i E_j}\right) - \frac{p_i \cdot p_j}{2} \left[\int_{-1}^1 \frac{dx}{p_x^2} \ln\left(\frac{k_{max}^2}{E_x^2}\right) + \int_{-1}^1 \frac{dx}{p_x^2} \ln\left(\frac{p_x^2}{\lambda^2}\right) \right] \right\}$$
Where,
$$2p_x = (1+x)p_i + (1-x)p_j$$

$$2E_x = (1+x)E_i + (1-x)E_j$$

$$2p'_x = (1+x)p_i \eta_i - (1-x)p_j$$







Virtual photon corrections

Due to contact term: Contain ultraviolet divergences

• Our method to construct the contact term provides O(e) term whereas the obtained UV divergence is at $\mathcal{O}(e^2)$.

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- Proposed solution: There may be higher dimensional operators to absorb this UV divergence or a new formalism is required to derive CT.
- Discarded the leftover UV divergences. The finite part is proportional to momenta of particles and numerically it contributes to $\sim 1.4\%$.





• Evaluating virtual diagrams: $M_{\text{virtual}} = M_0 \left[1 + M_0 \right]$

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The Coulomb factor :
$$\Omega_c = \prod_{i < j} \frac{-2\pi\alpha}{\beta_{ij}} \frac{1}{e^{\frac{-2\pi\alpha}{\beta_{ij}}} - 1} \qquad \text{where,} \qquad \beta_{ij} = \sqrt{1 - \frac{m_i^2 m_j^2}{(p_i \cdot p_j)^2}}$$

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where, $\beta_{ij} = \sqrt{1 - \frac{m_i^2 m_j^2}{(p_i \cdot p_j)^2}}$

$$K = \prod_{k \neq j} \frac{-2\pi\alpha}{\beta_{ij}} \frac{1}{e^{\frac{-2\pi\alpha}{\beta_{ij}}} - 1}$$
where, $\beta_{ij} = \sqrt{1 - \frac{m_i^2 m_j^2}{(p_i \cdot p_j)^2}}$





• The total decay rate: $d\Gamma_{real} = d\Gamma_0 \left(1 + 2\alpha \underbrace{(\tilde{\mathscr{B}} + \mathscr{B})}_{\mathscr{H}_{ij}} + \frac{\alpha}{\pi} \right) \Omega_c + d\Gamma'$







$$\alpha \underbrace{(\tilde{\mathscr{B}} + \mathscr{B})}_{\mathscr{H}_{ij}} + \frac{\alpha}{\pi} \Omega_c + d\Gamma'$$

$$\frac{1}{p_{j}} + \frac{p_{i} \cdot p_{j} \eta_{i} \eta_{j}}{2} \int_{-1}^{1} \frac{dx}{p_{x}^{'2}} \ln\left(\frac{p_{x}^{'2}}{\lambda^{2}}\right) + \frac{p_{i} \cdot p_{j}}{2} \int_{-1}^{1} \frac{dx}{p_{x}^{2}} \ln\left(\frac{k_{max}^{2} p_{x}^{2}}{E_{x}^{2} \lambda^{2}}\right)$$





• The correction factor $\Delta^{i}(\mathcal{O}(\alpha))$: $\Delta^{i} = \left(\frac{d^{2}\Gamma_{0}}{dsdq^{2}}\right)$



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• The correction factor $\Delta^i(\mathscr{O}(\alpha))$: $\Delta^i = \left(\frac{d^2 \Gamma_0}{ds dq^2}\right)^{-1} \left(\frac{d^2 \Gamma^i}{ds dq^2}\right) - 1$
• The shift $\Delta_{R_x^{ur}}$: $\Delta^i_{R_x^{ue}} = R_k^{0\mu e} \left(\frac{\Delta \Gamma^i_{\mu}}{\Gamma^i_{\mu}} - \frac{\Delta \Gamma^i_{e}}{\Gamma^i_{e}}\right)$













Results

• Correction factor for the electron is about three times larger than that for the muons (both are negative) and this difference is due to smallness of the electron mass compared to muon mass.





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Results

- Correction factor for the electron is about three times larger than that for the muons (both are negative) and this difference is due to smallness of the electron mass compared to muon mass.
- The QED corrections impact more massive charged particles significantly less compared to lighter particles.
- There is a mild dependence on the photon energy cut k_{max}
- It is sensitive to θ_{cut} , particularly for the case of electrons. Choosing $\theta_{cut} \sim$ few degrees, this sensitivity essentially disappears. For muon, it is not that sensitive.



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 We can see the explicit cancellation by choosing a different set of kinematical variables,

 $t = (p_B - p_k)^2$, $s = (p_k + p_2)^2$, $x = (p_k + k)^2$ and $q^2 = (p_2 + p_3)^2$ and E_k in the rest frame of $(q + k)^2$. With a different method results match with G. Isidori et. al.[JHEP 12 (2020) 104].



• We have fixed the contact term demanding the gauge invariance of the matrix amplitude.

Summary and Conclusions



Observ.	k_{max}	
	$(\theta_{cut}=3^\circ)$	(
Δ_e^c	125 MeV	(
Δ^c_{μ}	125 MeV	(
$R_{k^{\mu e}}$	125 MeV	(

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